

# Ultra-light Axion Dark Matter and CMB B-mode Polarization



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**Collaborators: Guo-Chin Liu (TKU)**

**Seokcheon Lee (NGU)**

# A HIGH-ENERGY NEUTRINO SIGNATURE FROM SUPERSYMMETRIC RELICS

John S. HAGELIN

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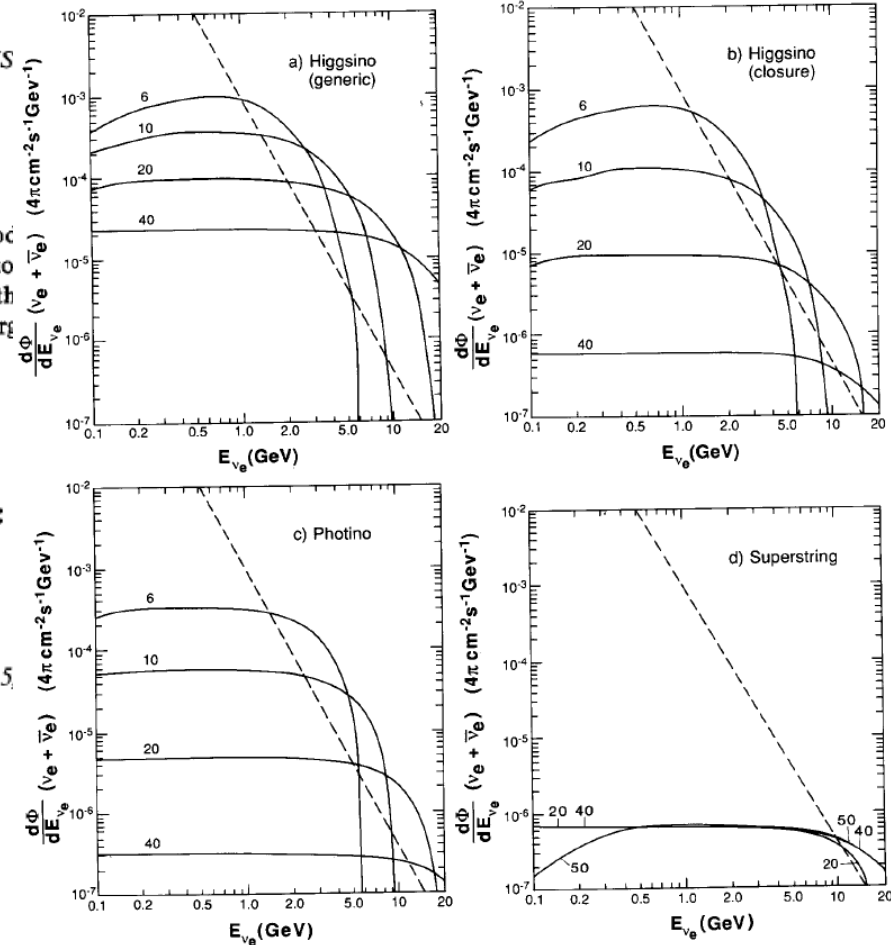
1987A Kamiokande, IMB,...

K.W. NG and Keith A. OLIVE

*School of Physics and Astronomy, University of Minnesota, Minneapolis, MN 55455, US*

Received 21 July 1986

We compute the energy spectrum of high-energy (0.1–10 GeV) neutrinos produced by the annihilation of supersymmetric (SUSY) cold dark matter trapped in the sun. We compare this spectrum to the solar flux of neutrinos and find that in the direction of the sun the solar flux of neutrinos can exceed the background at energies  $E_\nu \gtrsim 1$  GeV, and are as much as a factor  $\sim 30$  above background for energies  $E_\nu \gtrsim 10$  GeV for standard SUSY relics as well as for superstring relics.



## DARK MATTER INDUCED NEUTRINOS FROM THE SUN: THEORY VERSUS EXPERIMENT

Kin-Wang NG, Keith A. OLIVE

*School of Physics and Astronomy, University of Minnesota, Minneapolis, MN 55455*

and

Mark SREDNICKI

*Department of Physics, University of California, Santa Barbara, CA 93106, USA*

Received 30 December 1986

We analyse the restrictions from experimental data on high-energy neutrinos from the sun due to the annihilation of cold dark matter. We pay particular attention to uncertainties in the neutrino-, particle- and astro-physics. We conclude that photinos and higgsinos cannot be ruled out by the present IMB data. However, it appears safe to say that muon sneutrinos are not a viable candidate.

# Keith's small fuel-economy car in mid 80s





# Keiths' favorite food



# Dynamical Dark Energy and Its Coupling to Matter

Kin-Wang Ng



**Academia Sinica & LeCosPA, Taiwan &  
KIPAC, SLAC**

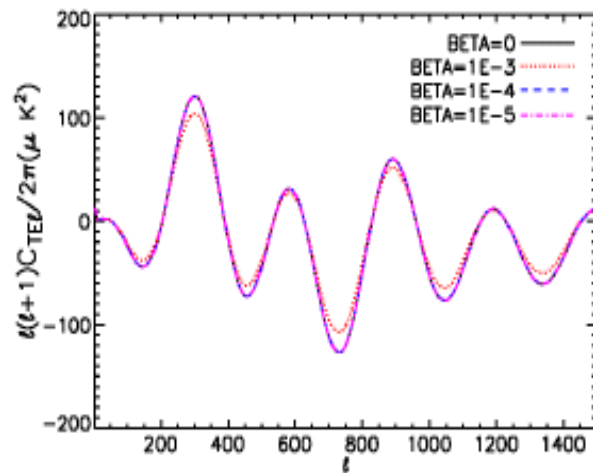
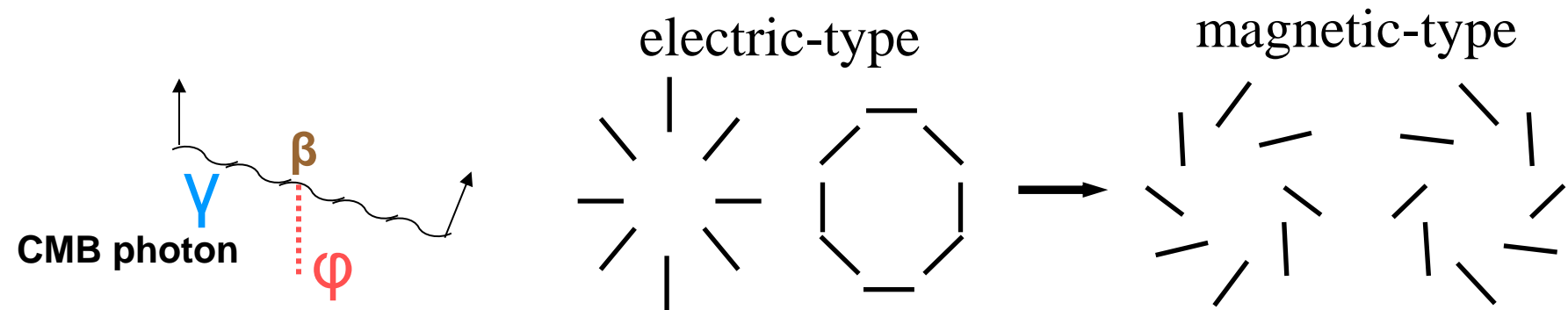


UM Cosmology Lunch Seminar Sep 28, 2009

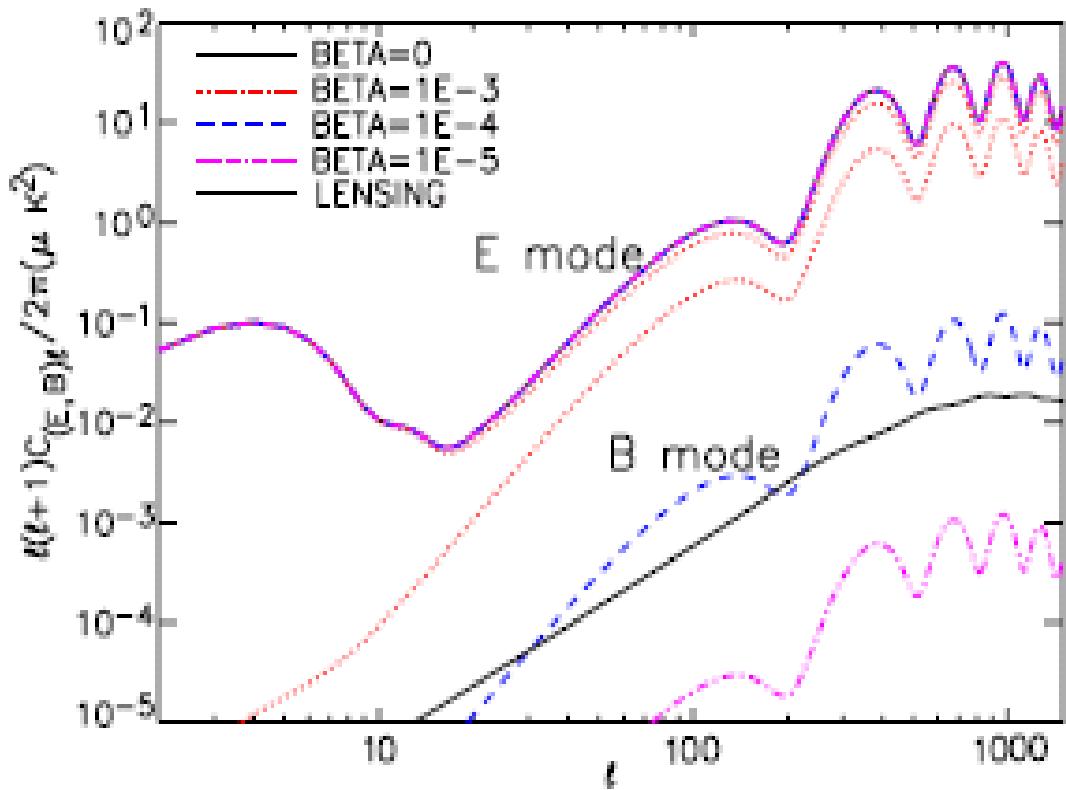
**Collaborators:** Da-Shin Lee (NDHU)  
Wolung Lee (NTNU)  
Guo-Chin Liu (TKU)  
Seokcheon Lee (ASloP)  
Chao-Lin Kuo (Stanford)

# DE mean field induced vacuum birefringence – cosmic rotation of CMB polarization

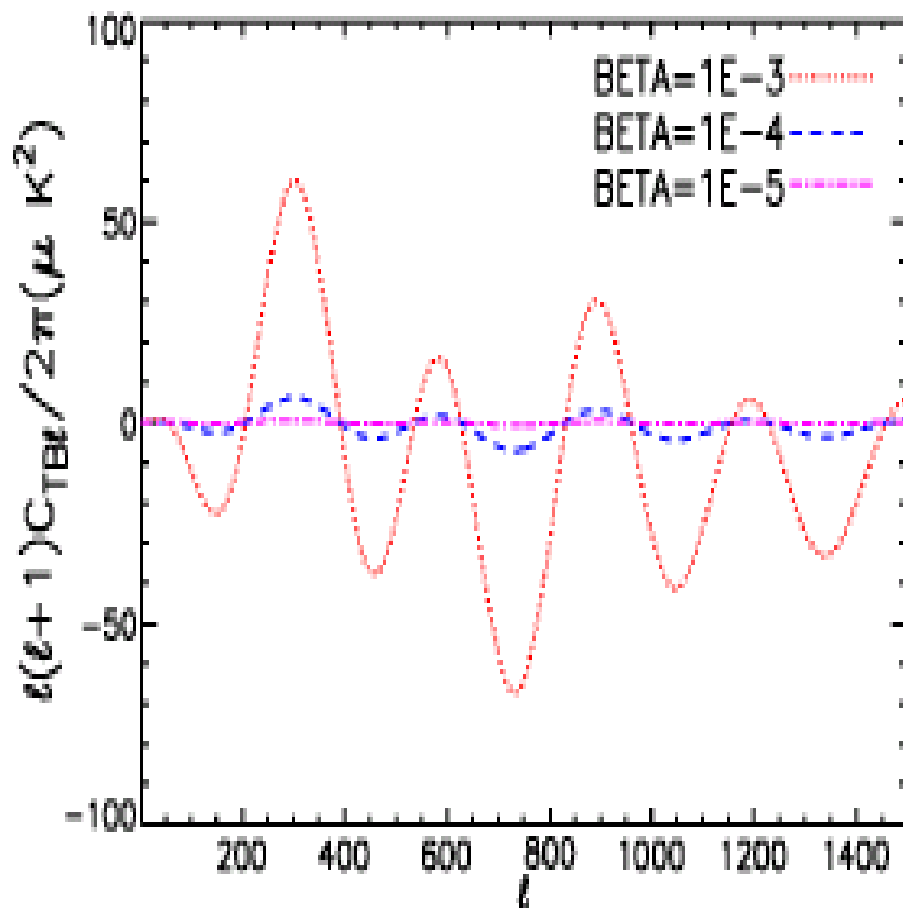
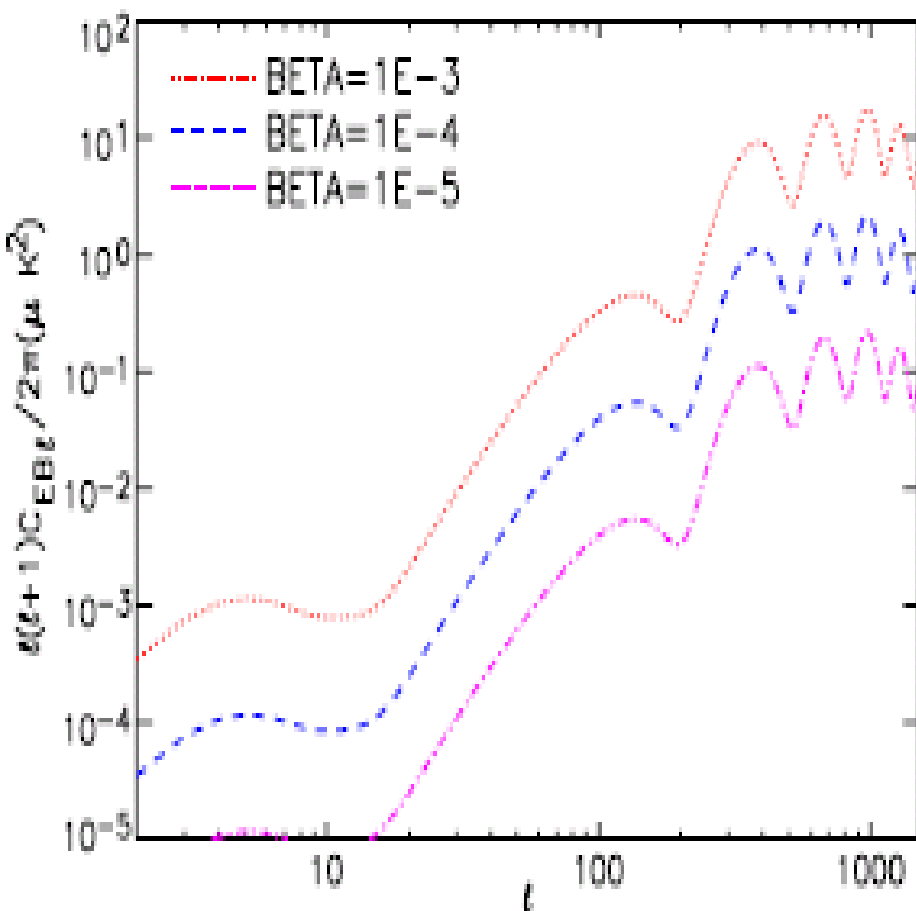
Liu, Lee, Ng 06



TE spectrum



# Parity violating EB,TB cross power spectra – cosmic parity violation



# Axion-like DE and CDM

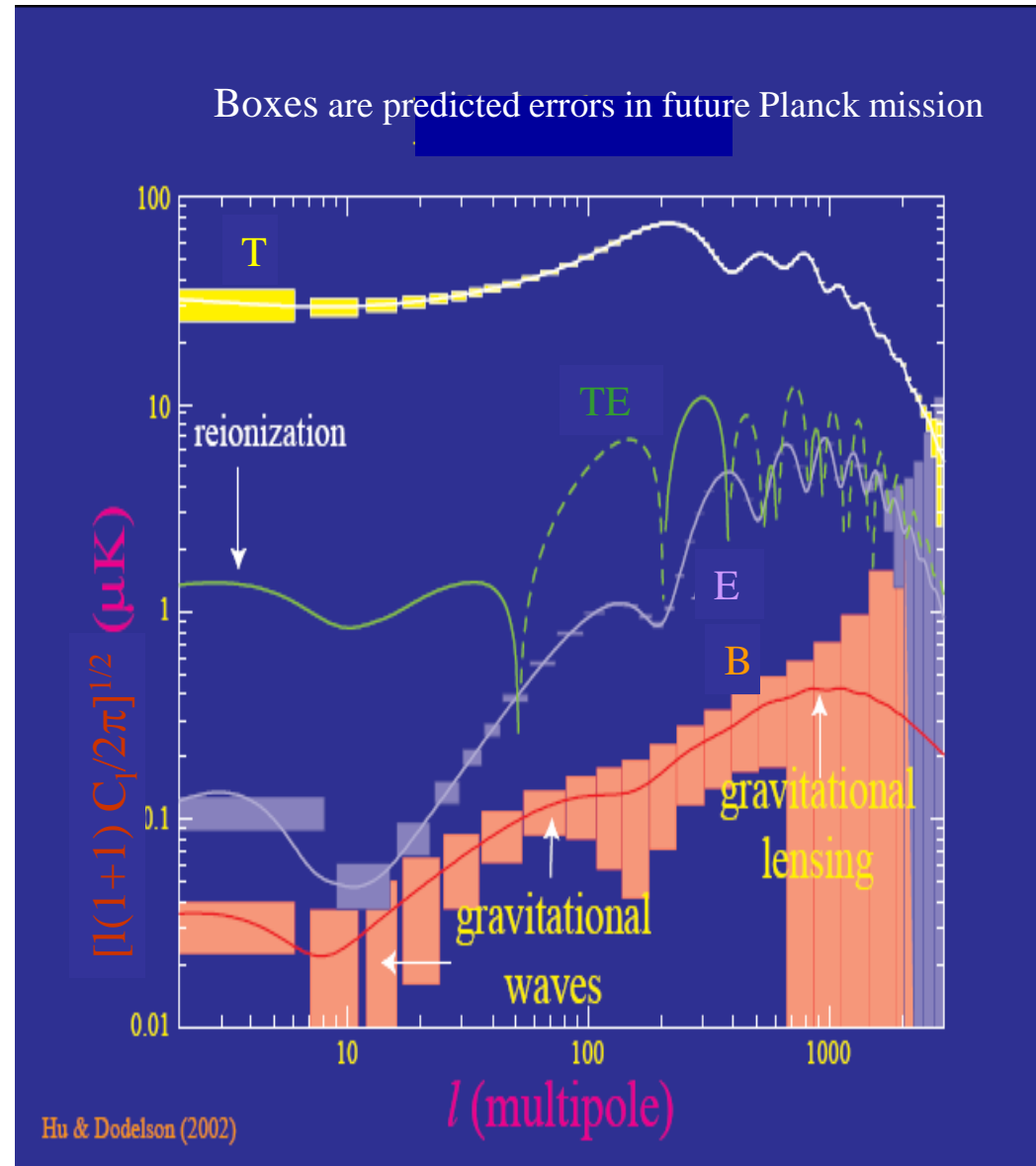
(too many references to list)

- Weak equivalence principle plus spin dictates a universal pseudoscalar (Ni 77)
- There exists at least one fundamental scalar – the Higgs boson !
- Axion monodromy – large-field inflation
- Peccei-Quinn symmetry breaking – QCD axion CDM
- Problems in small-scale structures –  $10^{-22}$  eV scalar (maybe pseudoscalar) fuzzy CDM
- String axiverse – a plentitude of axions with a vast mass range  $10^{-33}$  eV -  $10^{-10}$  eV Kaloper et al.
- Extended string axiverse – axions as DE

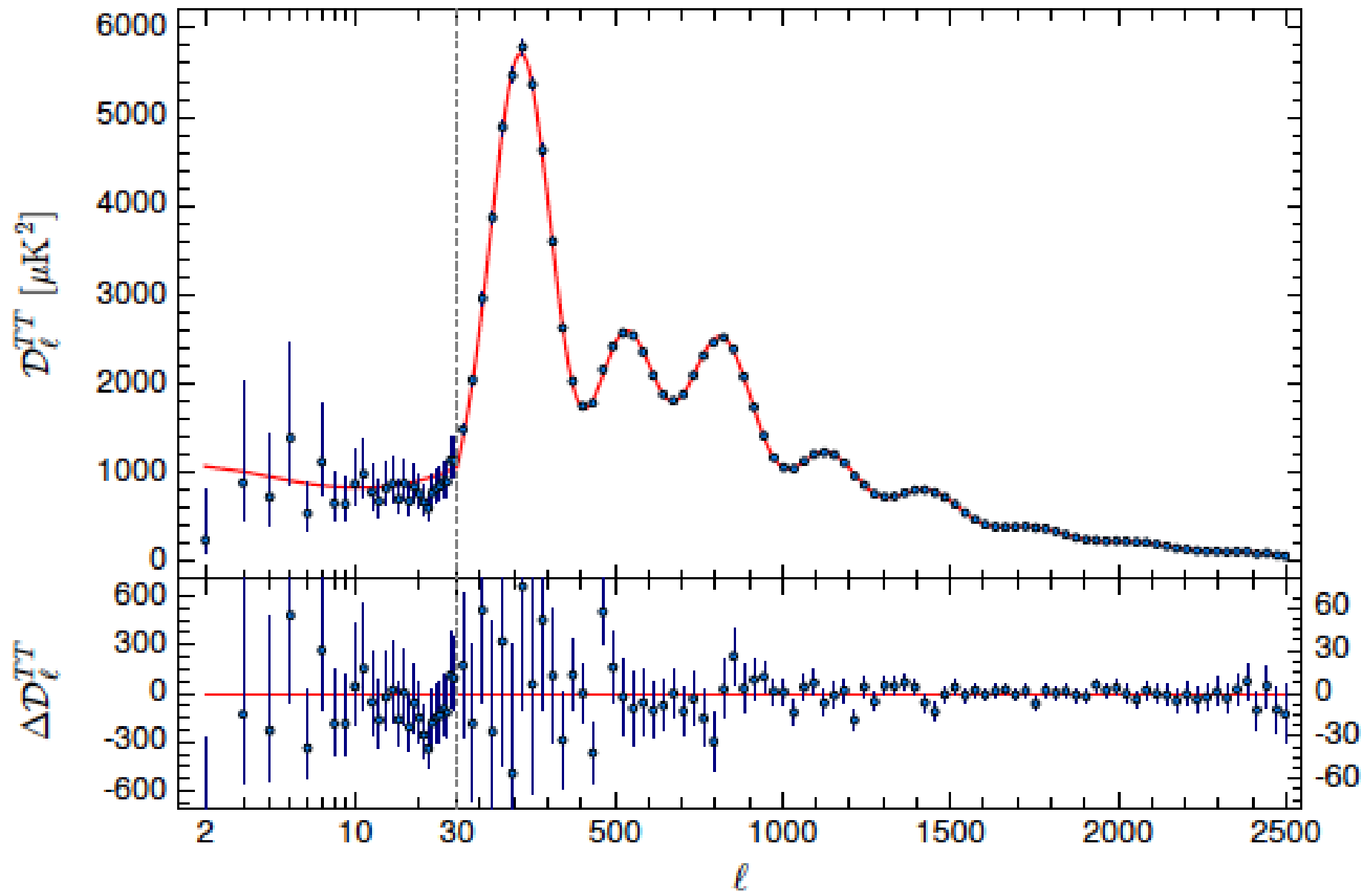


# Theoretical Predictions for CMB Power Spectra

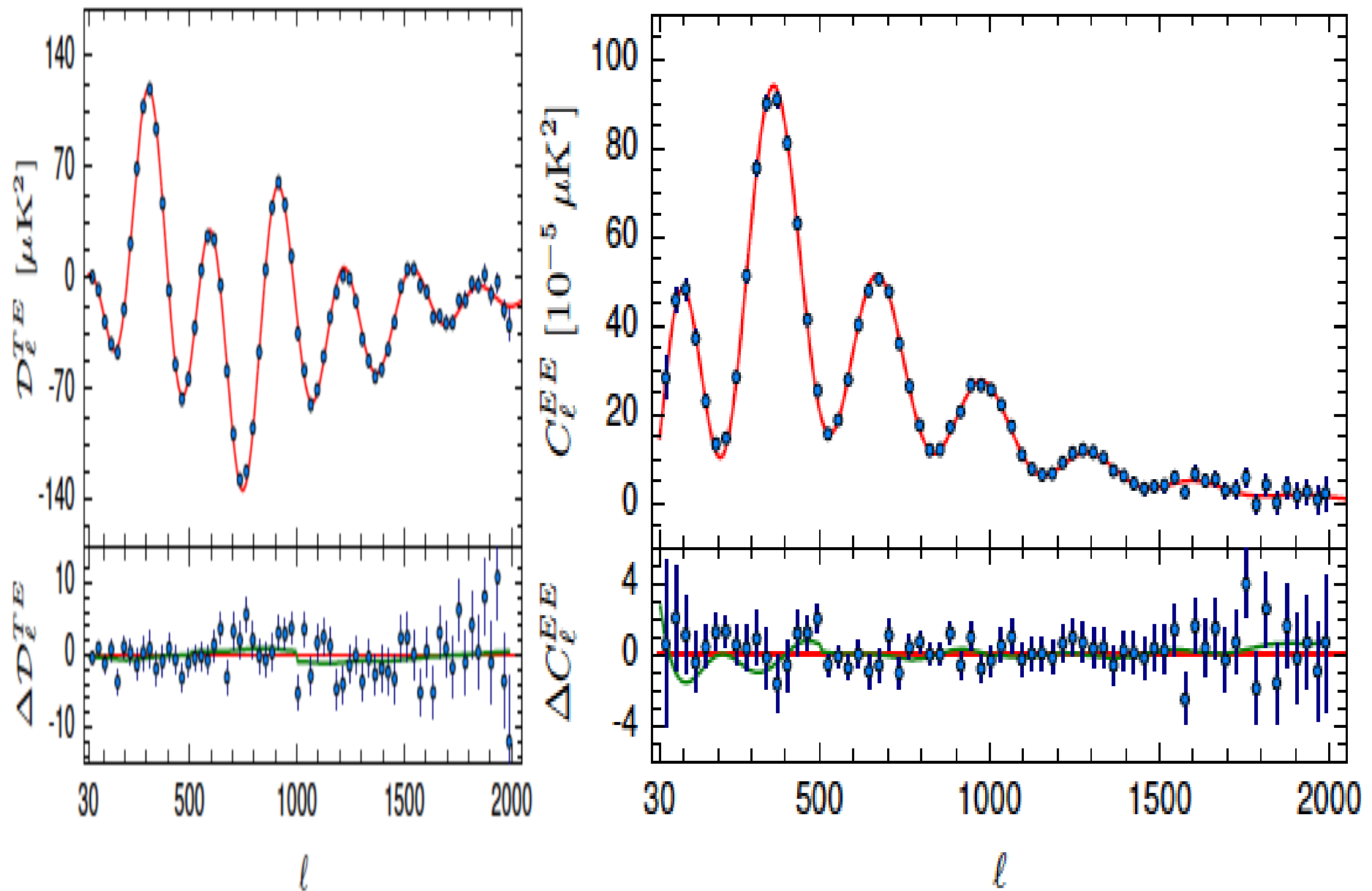
- Solving the radiative transfer equation for photons with electron scatterings
- Tracing the photons from the early ionized Universe through the last scattering surface to the present time
- Anisotropy induced by metric perturbations
- Polarization generated by photon-electron scatterings
- Power spectra dependent on the cosmic evolution governed by **cosmological parameters** such as matter content, density fluctuations, gravitational waves, ionization history, Hubble constant, and etc.



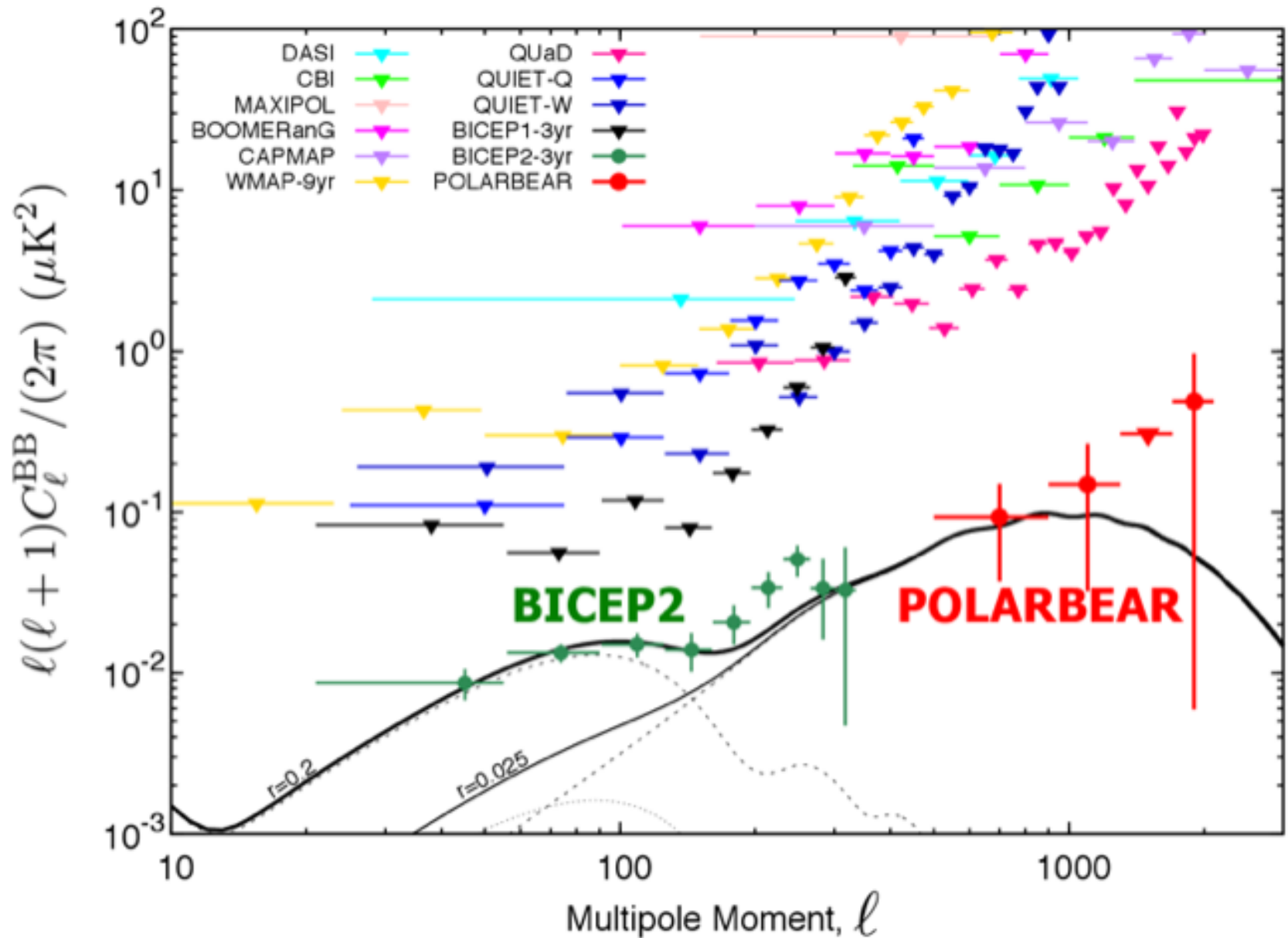
# Planck CMB Anisotropy $D_{\ell}^{TT} = 1(l+1) C_{\ell}^T$ 2015



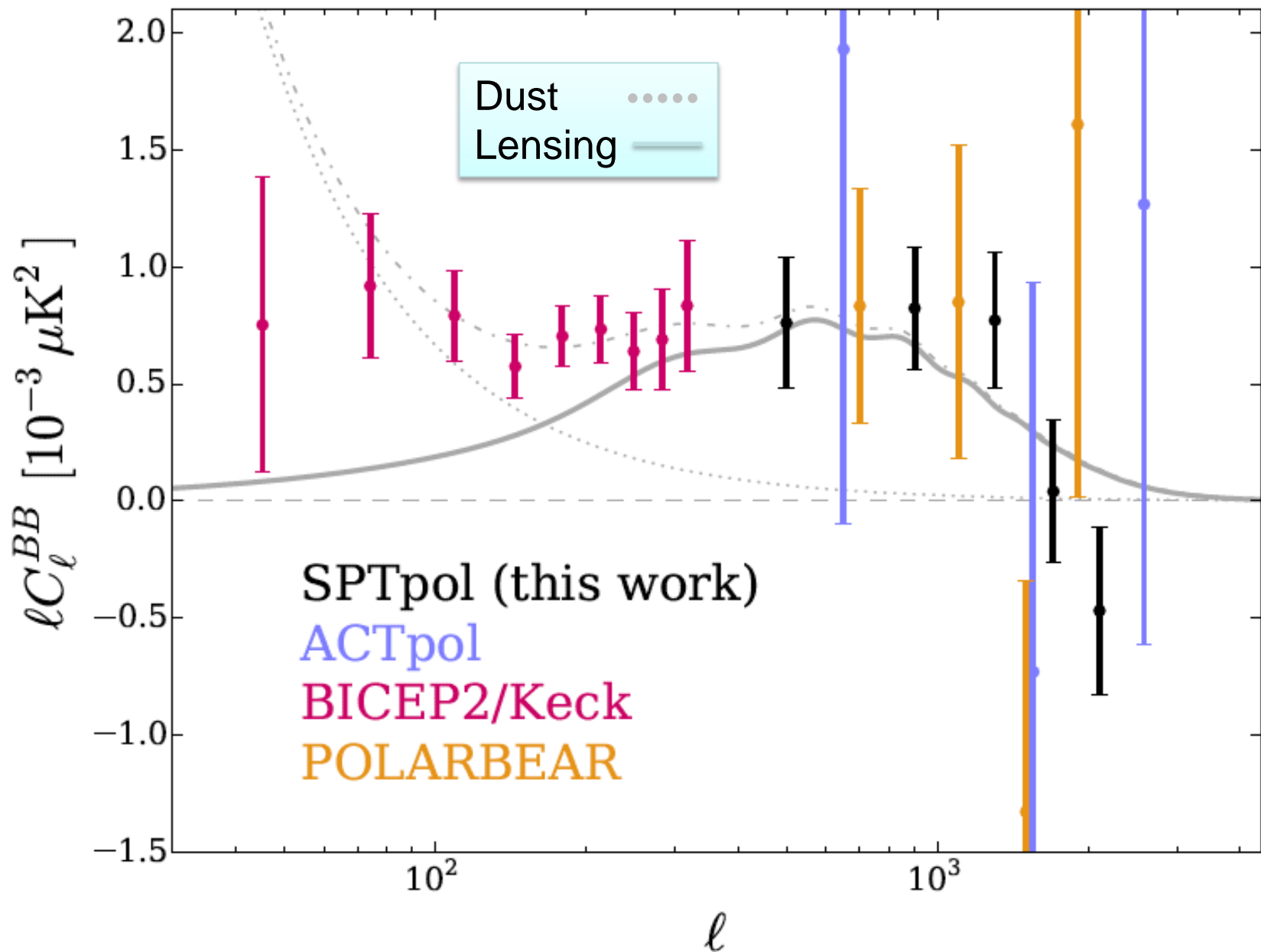
# Planck CMB Polarization Power Spectra 2015



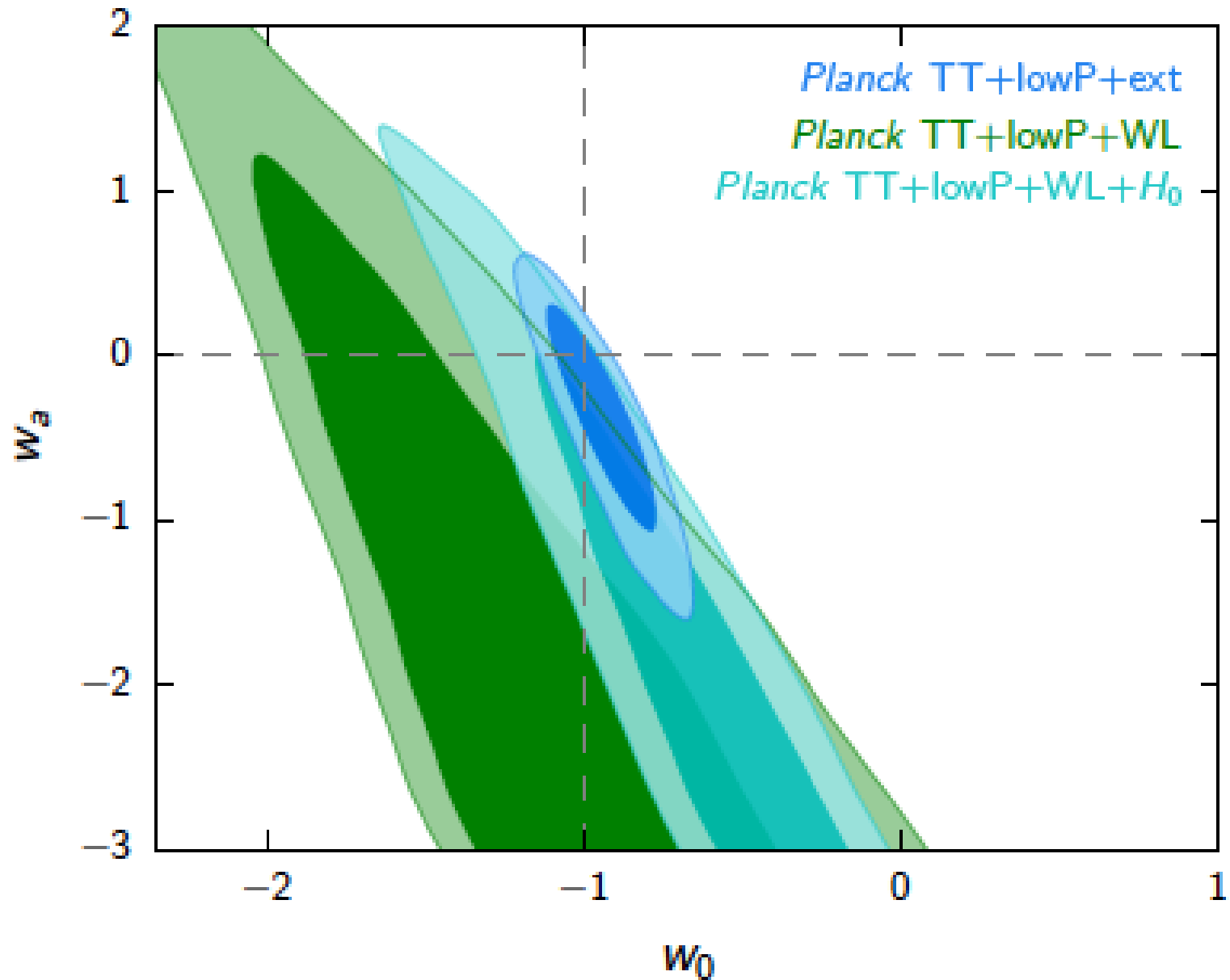
# POLARBEAR+BICEP2 B-mode Detection



# More B-mode detection



Time varying  $w=w_0+w_a z/(1+z)$





# DE/DM Coupling to Electromagnetism

$$\mathcal{L}_N = -\frac{1}{4}\sqrt{-g}B_{F\tilde{F}}(\phi)F_{\mu\nu}\tilde{F}^{\mu\nu}, \quad \text{where} \quad \phi \equiv \frac{\Phi}{M}, \quad M = M_{Pl}/\sqrt{8\pi}$$

This leads to photon dispersion relation Carroll, Field, Jackiw 90

$$n_{\pm} = \varepsilon \mp \frac{1}{2} \frac{\partial B_{F\tilde{F}}}{\partial \phi} \left( \frac{\partial \phi}{\partial \eta} + \vec{\nabla} \phi \cdot \hat{n} \right) \quad (\varepsilon, \vec{n}) \text{ is the photon four-momentum}$$

$\pm$  left/right handed     $\eta$  conformal time

**vacuum birefringence**

then, a rotational speed of polarization plane

$$\omega = \frac{1}{2}(n_+ - n_-) = -\frac{1}{2} \frac{\partial B_{F\tilde{F}}}{\partial \phi} \left( \frac{\partial \phi}{\partial \eta} + \vec{\nabla} \phi \cdot \hat{n} \right)$$

If  $B=\beta\phi$ , cooling of horizontal branch stars would imply  $\beta < 10^7$

# Radiative transfer equation

$$\dot{\Delta}_{Q\pm iU}(\vec{k}, \eta) + ik\mu\Delta_{Q\pm iU}(\vec{k}, \eta) = n_e\sigma_T a(\eta) \left[ -\Delta_{Q\pm iU}(\vec{k}, \eta) \right. \\ \left. \sum_m \sqrt{\frac{6\pi}{5}} {}_{\pm 2}Y_2^m(\hat{n}) S_P^{(m)}(\vec{k}, \eta) \right] \mp i2\omega\Delta_{Q\mp iU}(\vec{k}, \eta),$$

$\mu = \mathbf{n} \cdot \mathbf{k}$ ,  
 $\eta$ : conformal time  
 $a$ : scale factor  
 $n_e$ : e density  
 $\sigma_T$ : Thomson cross section

**Source term for polarization**

$$S_P^{(m)}(\vec{k}, \eta) \equiv \Delta_{T_2}^{(m)}(\vec{k}, \eta) + 12\sqrt{6}\Delta_{+,2}^{(m)}(\vec{k}, \eta) + 12\sqrt{6}\Delta_{-,2}^{(m)}(\vec{k}, \eta)$$

**Cosmic rotation**

$$\bar{\omega}(\eta) = -\frac{1}{2}\beta_{F\tilde{F}} \frac{d\bar{\phi}}{d\eta}.$$

$$\mathcal{L}_N = -\frac{1}{4}\sqrt{-g}B_{F\tilde{F}}(\phi)F_{\mu\nu}\tilde{F}^{\mu\nu}$$

$$B_{F\tilde{F}} = \beta_{F\tilde{F}}\phi \quad \phi \equiv \frac{\Phi}{M}$$

**Rotation angle**

$$\alpha(\eta) = \int_{\eta}^{\eta^0} d\eta' \omega(\eta')$$

# Power spectra

$$C_{\ell}^{(E,B)} = (4\pi)^2 \frac{9}{16} \frac{(\ell+2)!}{(\ell-2)!} \int k^2 dk \left[ \Delta_{(E,B)}(k, \eta_0) \right]^2,$$

$$C_{\ell}^{EB} = (4\pi)^2 \frac{9}{16} \frac{(\ell+2)!}{(\ell-2)!} \int k^2 dk \Delta_E(k, \eta_0) \Delta_B(k, \eta_0),$$

$$C_{\ell}^{TE} = (4\pi)^2 \sqrt{\frac{9}{16} \frac{(\ell+2)!}{(\ell-2)!}} \int k^2 dk \Delta_T(k, \eta_0) \Delta_E(k, \eta_0),$$

$$C_{\ell}^{TB} = (4\pi)^2 \sqrt{\frac{9}{16} \frac{(\ell+2)!}{(\ell-2)!}} \int k^2 dk \Delta_T(k, \eta_0) \Delta_B(k, \eta_0),$$

$$\Delta_T(k, \eta_0) = \int_0^{\eta_0} d\eta g(\eta) S_T(k, \eta) j_{\ell}(kr),$$

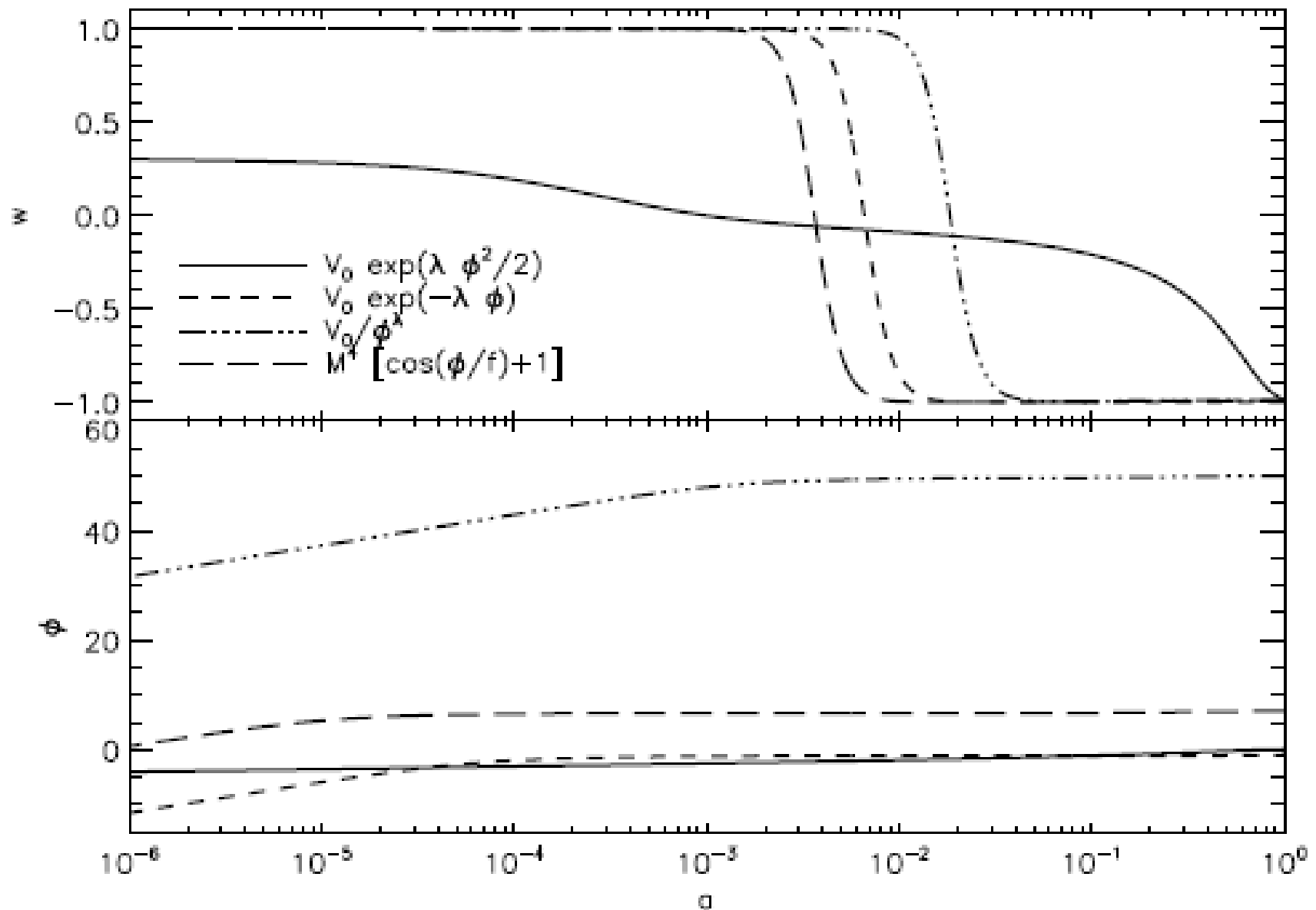
$$\Delta_E(k, \eta_0) + i\Delta_B(k, \eta_0) = \int_0^{\eta_0} d\eta g(\eta) S_P(k, \eta) \frac{j_{\ell}(kr)}{(kr)^2} e^{i2\alpha(\eta)},$$

**$g(\eta)$ : radiative transfer function**

**$S_T$ : source term for anisotropy**

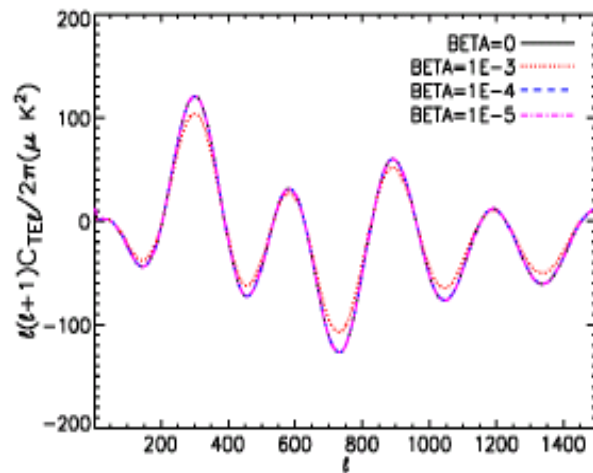
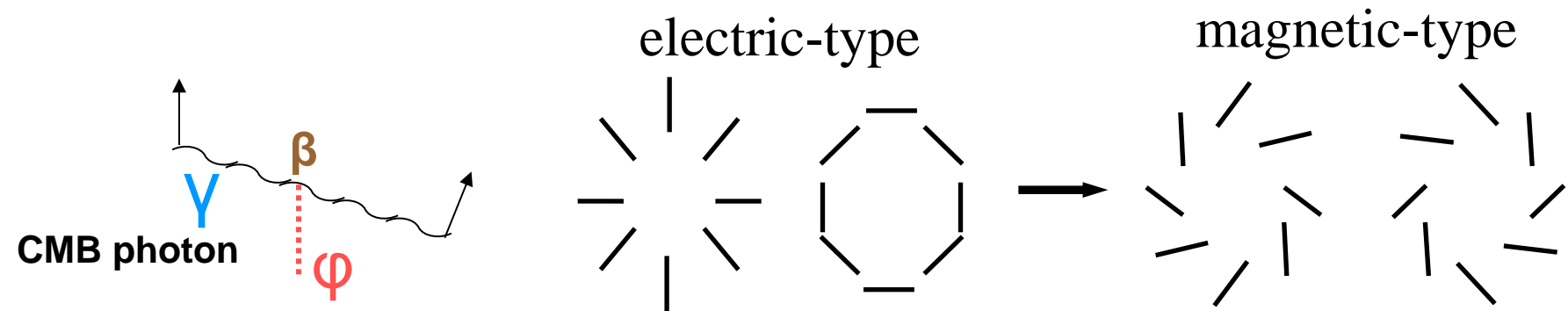
**$S_P = S_P^{(0)}$        $r = \eta_0 - \eta$**

# We Tried Many Scalar Dark Energy Models

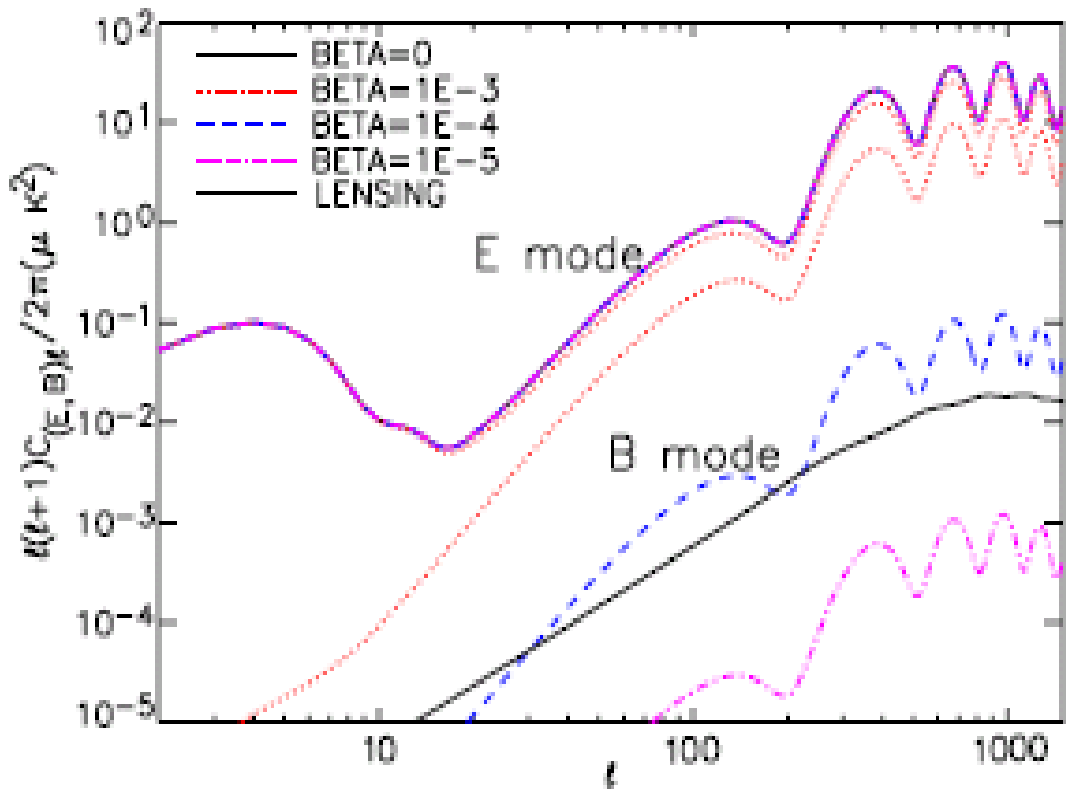


# DE mean field induced vacuum birefringence – cosmic rotation of CMB polarization

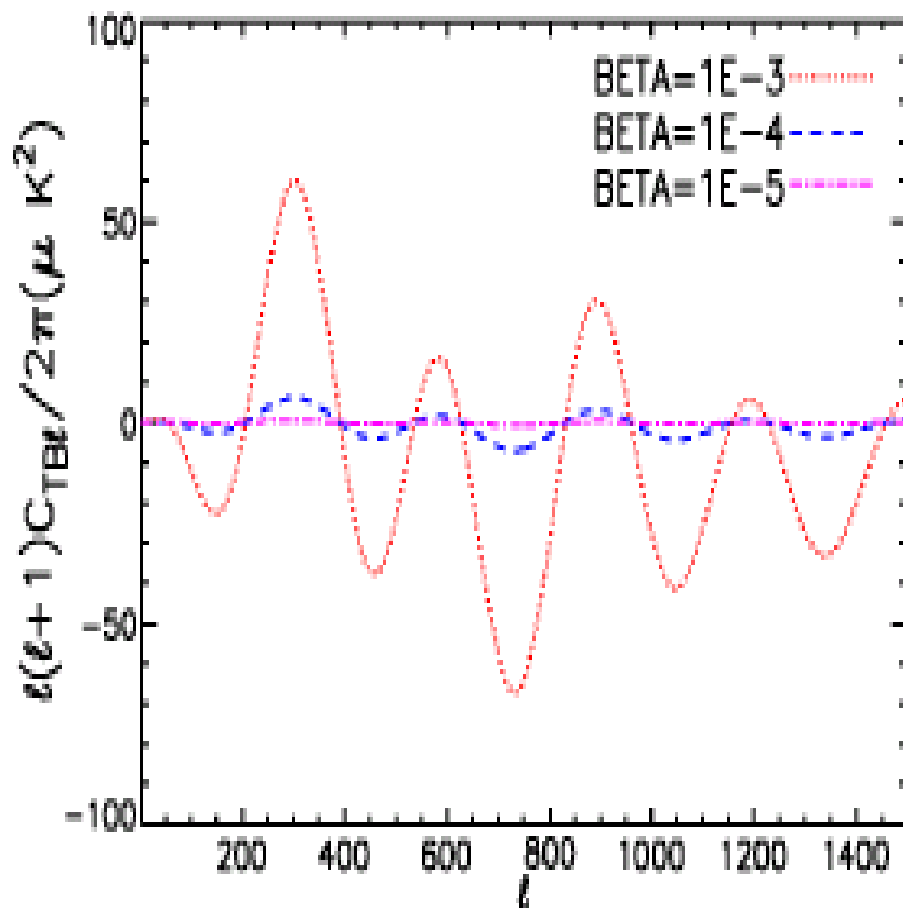
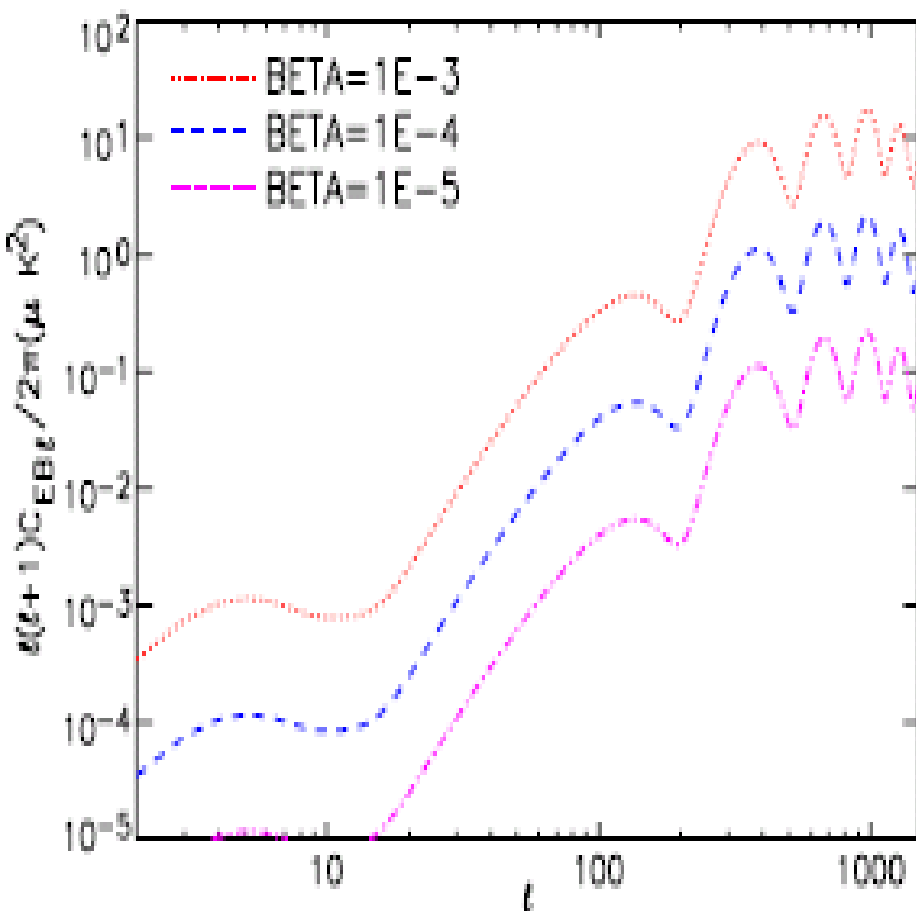
Liu, Lee, Ng 06



TE spectrum



# Parity violating EB,TB cross power spectra – cosmic parity violation





# Including Dark Energy Perturbation

Dark energy  
perturbation

$$\phi(\eta, \vec{x}) = \bar{\phi}(\eta) + \delta\phi(\eta, \vec{x}) \quad \delta\phi(\eta, \vec{x}) = \frac{1}{\sqrt{(2\pi)^3}} \int \delta\phi(\vec{k}', \eta) e^{i\vec{k}' \cdot \vec{x}} d^3k'$$

time and space  
dependent rotation

$$\omega = -\frac{1}{2} \frac{\partial B_{F\tilde{F}}}{\partial \phi} \left( \frac{\partial \phi}{\partial \eta} + \vec{\nabla} \phi \cdot \hat{n} \right)$$

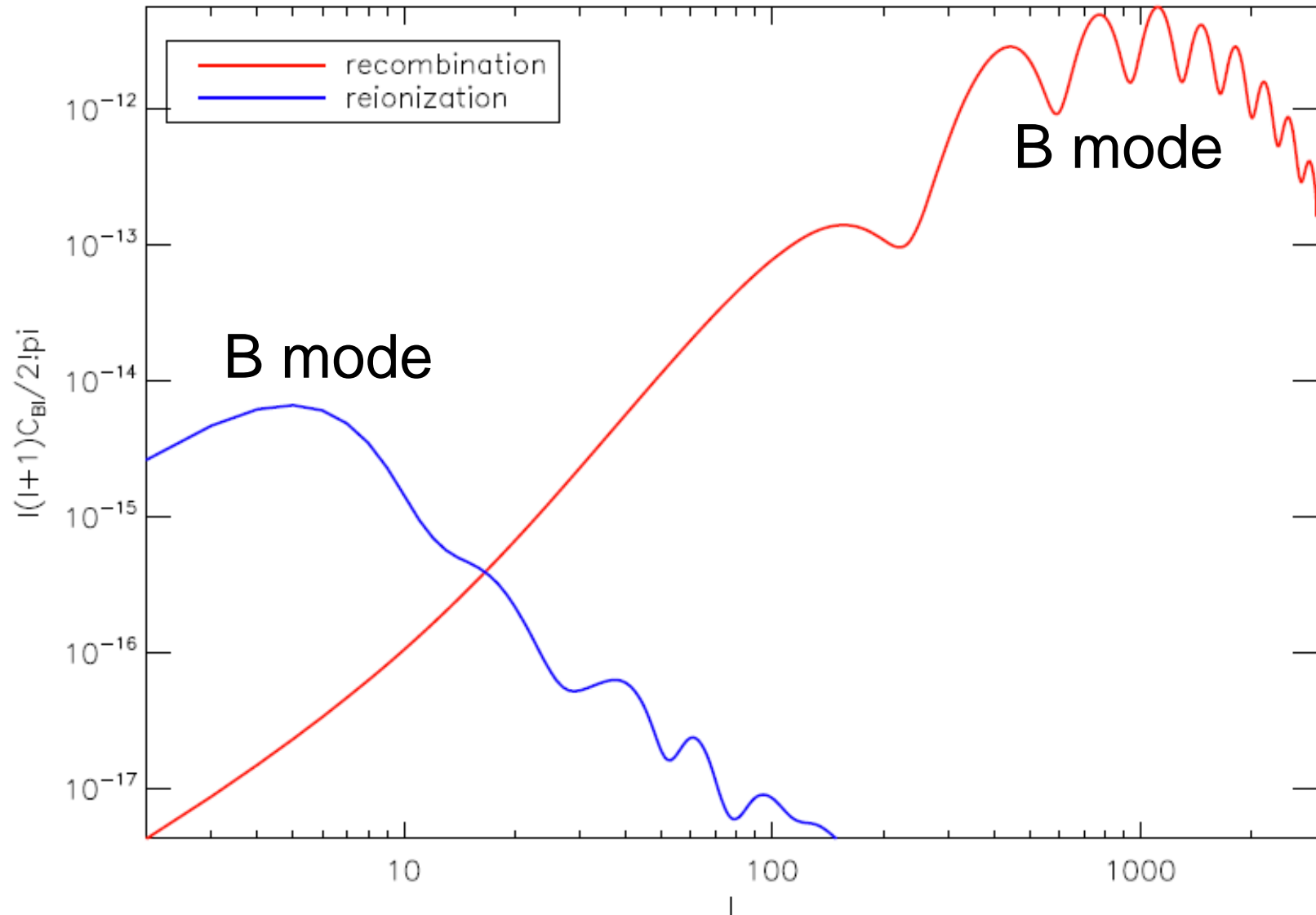
$$\dot{\Delta}_{Q\pm iU}(\vec{k}, \eta) + ik\mu\Delta_{Q\pm iU}(\vec{k}, \eta) = n_e\sigma_T a(\eta) \left[ -\Delta_{Q\pm iU}(\vec{k}, \eta) \times \sum_m \sqrt{\frac{6\pi}{5}} {}_{\pm 2}Y_2^m(\hat{n}) S_P^{(m)}(\vec{k}, \eta) \right] \mp i2 \frac{1}{\sqrt{(2\pi)^3}} \int d\vec{k}' \tilde{\omega}(\vec{k} - \vec{k}', \eta) \Delta_{Q\pm iU}(\vec{k}', \eta)$$

$$\tilde{\omega}(\vec{k}, \eta) = -\frac{1}{2} \frac{\partial B_{F\tilde{F}}}{\partial \bar{\phi}} \left[ \dot{\delta\phi}_{\vec{k}}(\eta) + i\vec{k} \cdot \hat{n} \delta\phi_{\vec{k}}(\eta) \right]$$

- Perturbation induced polarization power spectra in general quintessence models are small
- Interestingly, in nearly  $\Lambda$ CDM models (no time evolution of the mean field), birefringence generates  $\langle BB \rangle$  while  $\langle TB \rangle = \langle EB \rangle = 0$

Dark energy perturbation with  $w=-1$  Lee,Liu,Ng 14

Birefringence generates  $\langle BB \rangle$  while  $\langle TB \rangle = \langle EB \rangle = 0$



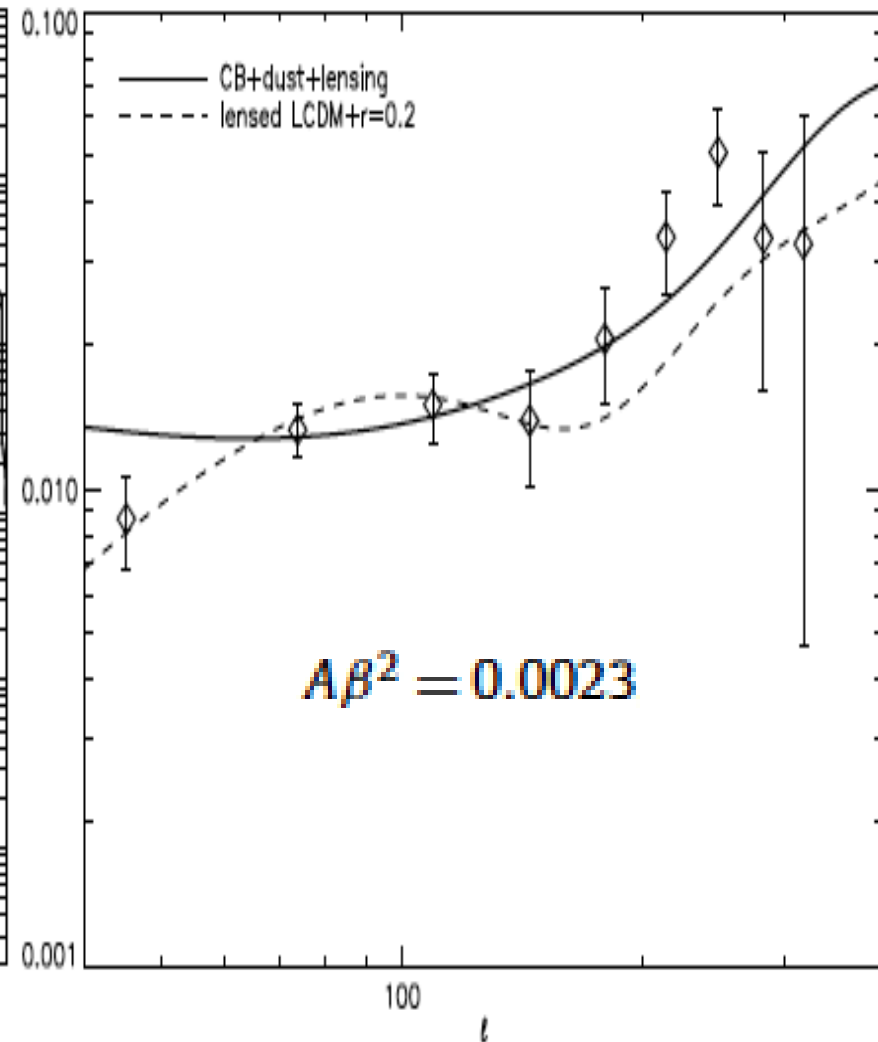
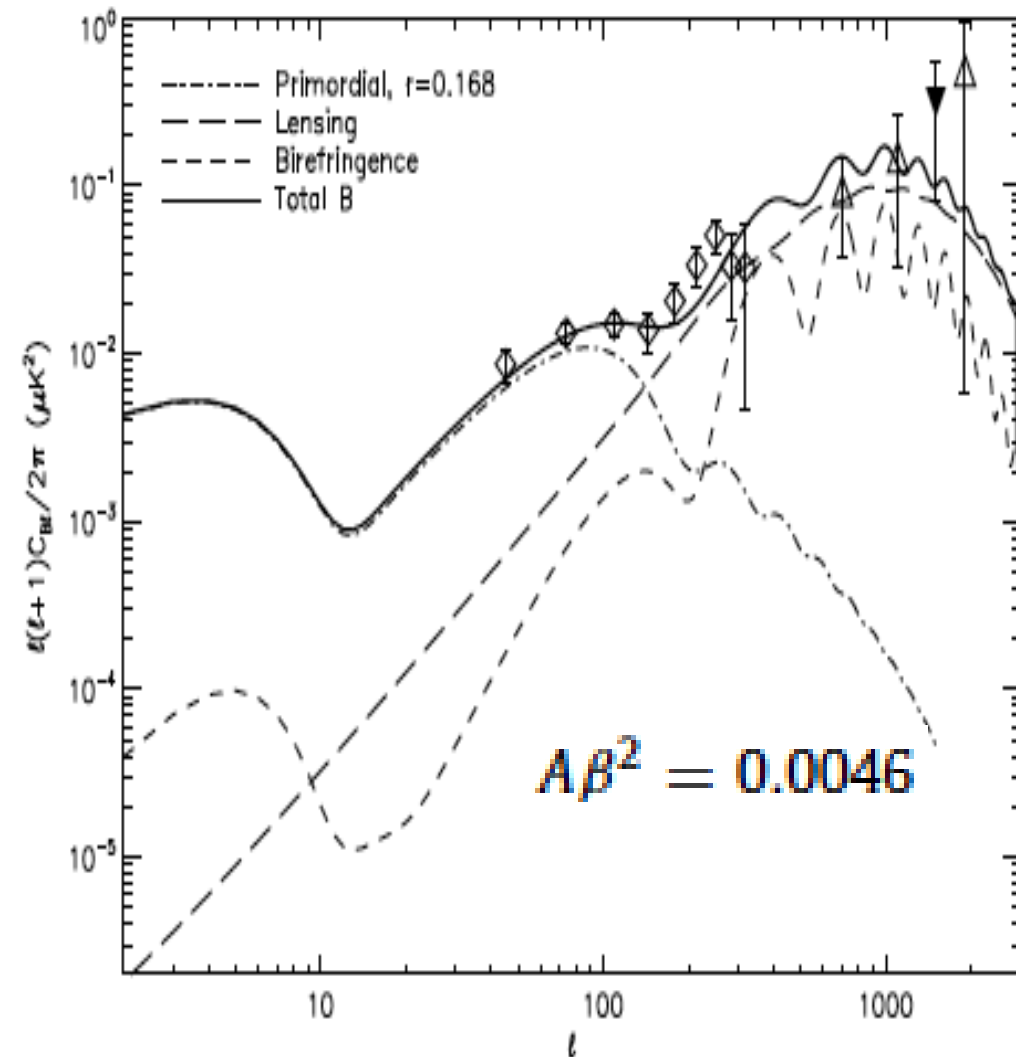
# Cosmic Birefringence (CB) Fluctuations

Nearly massless  
pseudo scalar

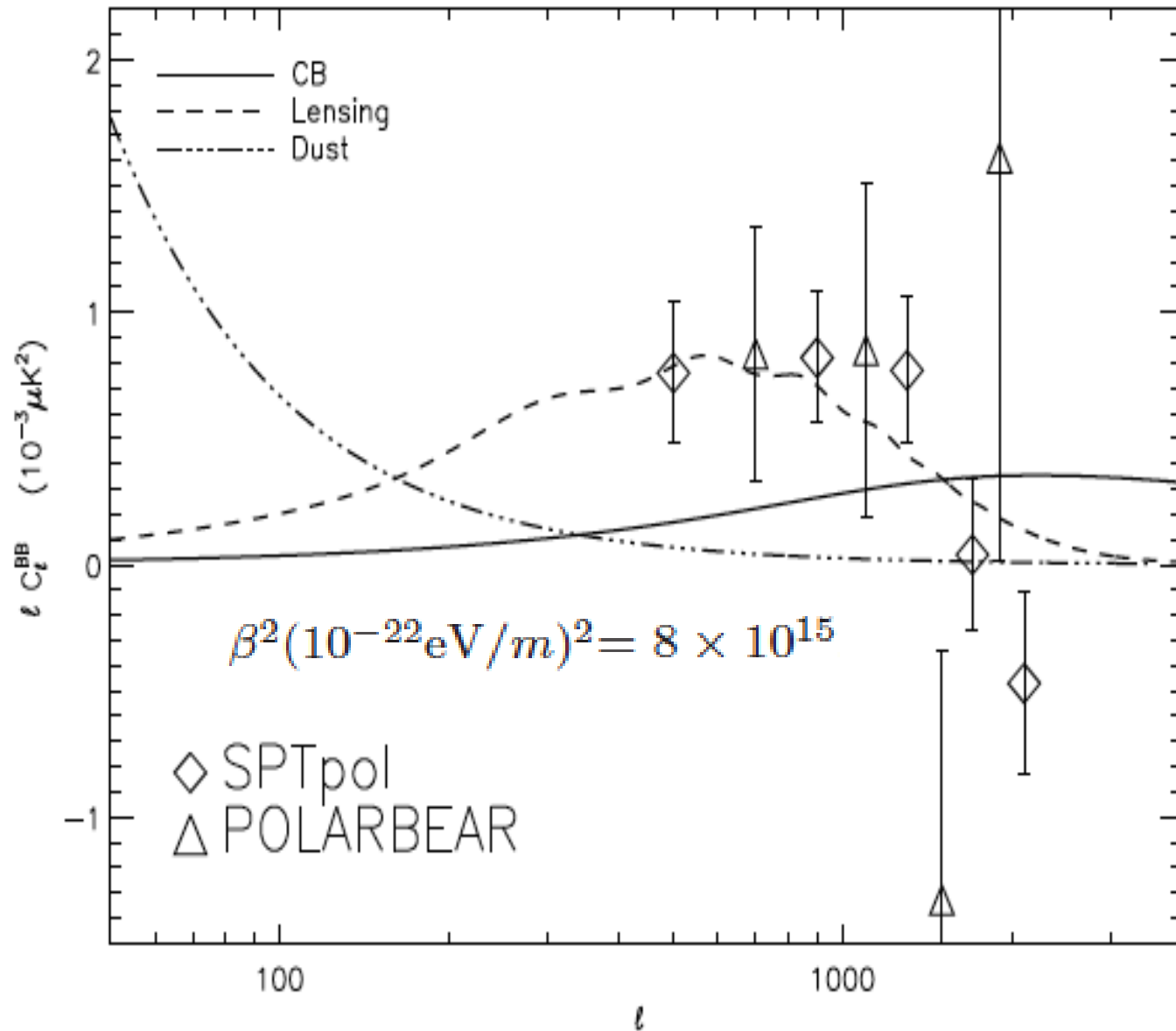
$$\langle \delta\phi_{\vec{k},i} \delta\phi_{\vec{k}',i} \rangle = (2\pi^2/k^3) P_{\delta\phi}(k) \delta(\vec{k} - \vec{k}')$$

$$P_{\delta\phi}(k) = A k^{n-1}$$

Pospelov et al 09  
Lee, Liu, Ng 14



# Axion ( $m \sim 10^{-22} \text{eV}$ ) CDM curvature perturbation



$$\delta \rho_\phi / \rho_\phi = \delta \rho / \rho$$

$$\rho_\phi = m_\phi^2 \phi^2$$

$$\frac{\delta \rho_\phi}{\rho_\phi} = 2 \frac{\delta \phi}{\phi}$$

$$\frac{\delta \phi}{M} = \frac{\delta \phi}{\phi} \frac{\phi}{M}$$

$$\phi = \phi_m \left( \frac{a_m}{a} \right)^{\frac{3}{2}}$$

$$H(a_m) = m$$

# Summary

- Using CMB B-mode polarization to search for dark energy induced **vacuum birefringence**
  - Mean field time evolution  $\rightarrow \langle BB \rangle, \langle TB \rangle, \langle EB \rangle$
  - Include DE perturbation  $\rightarrow \langle BB \rangle, \langle TB \rangle = \langle EB \rangle = 0$
- Axion cold matter matter curvature perturbation  $\rightarrow \langle BB \rangle, \langle TB \rangle = \langle EB \rangle = 0$ ; isocurvature perturbation?
- This may confuse the searching for genuine B modes induced by gravitational lensing or primordial gravitational waves, so de-rotation is needed to remove vacuum birefringence effects [Kamionkowski 09](#), [Ng 10](#)